Reconstruction of time-dependent sources and scalar parameters of fractional diffusion equations from final measurements

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Notation of fractional integral derivatives and Riemann-Liouville and Caputo fractional derivatives:

$$I^{\beta}w(t) = \int_0^t \frac{(t-\tau)^{\beta-1}}{\Gamma(\beta)} w(\tau) d\tau,$$

$${}^RD^{\beta}w(t) = \frac{d}{dt} \int_0^t \frac{(t-\tau)^{-\beta}}{\Gamma(1-\beta)} w(\tau) d\tau,$$

$${}^CD^{\beta}w(t) = \int_0^t \frac{(t-\tau)^{-\beta}}{\Gamma(1-\beta)} w'(\tau) d\tau,$$

where $0 < \beta < 1$.

Fractional diffusion equation in Riemann-Liouville form:

$$u_t(t,x) = {}^R D^{1-\beta}(\lambda \Delta u)(t,x) + Q(t,x),$$

where $\lambda > 0$, Δ is the Laplacian and Q is the source term.

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The same equation in Caputo form:

$$^{C}D^{\beta}u(t,x) = \lambda \Delta u(t,x) + F(t,x),$$
 where $F = I^{1-\beta}Q$.

Let Q have the form

$$Q(t,x)=q(t)f(x).$$

Then

$$F(t,x)=I^{1-\beta}q(t)f(x).$$

Let
$$\Omega \subset \mathbb{R}^m$$
 be open bounded domain. Direct problem: ${}^CD^{\beta}u(t,x) = \lambda \Delta u(t,x) + I^{1-\beta}q(t)f(x), \ t \in (0,T), \ x \in \Omega, \ u(t,x) = 0, \ t \in (0,T), \ x \in \partial \Omega, \ u(0,x) = \varphi(x), \ x \in \Omega.$

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Final overdetermination condition:

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Eigenvalues and eigenfunctions of $-\Delta$:

$$-\Delta v_k(x) = \mu_k v_k(x), \ x \in \Omega, \quad v_k(x) = 0, \ x \in \partial \Omega,$$

It holds $0 < \mu_1 \le \mu_2 \le \ldots$, $\mu_k \to \infty$ and ν_k , $k \in \mathbb{N}$ constitute a complete orthonormal system in $L_2(\Omega)$.

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Eigenvalues and eigenfunctions of $-\lambda \Delta$:

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It holds

$$\lambda_k = \lambda \mu_k.$$



Expansions of data and state function u:

$$f(x) = \sum_{k=1}^{\infty} f_k v_k(x),$$

$$\varphi(x) = \sum_{k=1}^{\infty} \varphi_k v_k(x),$$

$$\psi(x) = \sum_{k=1}^{\infty} \psi_k v_k(x),$$

$$u(t, x) = \sum_{k=1}^{\infty} u_k(t) v_k(x).$$

$$u_k(t) = \varphi_k E_{\beta}(-\lambda_k t^{\beta}) + f_k \int_0^t (t - \tau)^{\beta - 1} E_{\beta, \beta}(-\lambda_k (t - \tau)^{\beta}) I^{1 - \beta} q(\tau) d\tau,$$

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into account we obtain the following basic relations:

$$\psi_{k} = \varphi_{k} E_{\beta}(-\lambda_{k} T^{\beta}) + f_{k} \int_{0}^{T} E_{\beta}(-\lambda_{k} (T - \tau)^{\beta}) q(\tau) d\tau, \quad k \in \mathbb{N}. \quad (3)$$

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Then

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Next we introduce the following restriction for q:

$$\exists \delta \in (0,T) : q(t) = 0, t \in (T - \delta, T).$$

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We use the asymptotics of Mittag-Leffler function

$$E_{eta}(-z) = -\sum_{n=1}^N rac{(-1)^n}{\Gamma(1-neta)z^n} + O\left(rac{1}{z^{N+1}}
ight) \quad ext{as} \quad z o \infty, \quad N \in \mathbb{N}.$$

$$0 = -\sum_{n=1}^{N} \frac{(-1)^n}{\lambda_{k_i}^n} \frac{1}{\Gamma(1 - n\beta)} \int_0^{T - \delta} \frac{q(\tau)}{(T - \tau)^{n\beta}} d\tau + O\left(\frac{1}{\lambda_{k_i}^{N+1}}\right)$$
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as $i \to \infty$, $N \in \mathbb{N}$.

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Multiplying (7) by λ_{k_i} and passing to the limit $i \to \infty$ we obtain

$$rac{1}{\Gamma(1-eta)}\int_0^{T-\delta}rac{q(au)}{(T- au)^eta}d au=0.$$

This means that the 1st addend under the sum in (7) is 0.

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Continuing this process we deduce the relations

$$\frac{1}{\Gamma(1-n\beta)}\int_0^{T-\delta}\frac{q(\tau)}{(T-\tau)^{n\beta}}d\tau=0, \quad n\in\mathbb{N}.$$
 (8)

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may not be useful.

Since the gamma function has poles at nonpositive integers, it holds

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Therefore, we consider separately two cases:

$$\beta \in (0,1) \setminus \mathbb{Q}$$
 and $\beta \in (0,1) \cap \mathbb{Q}$.

Firstly, let $\beta \in (0,1) \setminus \mathbb{Q}$. Then $\frac{1}{\Gamma(1-n\beta)} \neq 0$, $n \in \mathbb{N}$, hence

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Therefore,

$$\int_{\frac{1}{\tau\beta}}^{\frac{1}{\delta\beta}} s^n v(s) ds = 0, \quad n \in \{0\} \cup \mathbb{N}, \quad v(s) = s^{-\frac{1}{\beta}} q(T - s^{-\frac{1}{\beta}}).$$

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Set of polynomials is dense in $L_2(\frac{1}{T^{\beta}}, \frac{1}{\delta^{\beta}})$. Therefore,

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So we pick up the basic relations with such n:

$$\frac{1}{\Gamma(1-n'l_1-\beta)}\int_0^{T-\delta}\frac{q(\tau)}{(T-\tau)^{n'l_1+\beta}}d\tau=0,\quad n'\in\mathbb{N}.$$



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The proof can be finished as in the previous case



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Sufficient condition for (9):

If
$$\varphi \in L_2(\Omega) \setminus D((-L)^{\alpha})$$
 and $f \in D((-L)^{\alpha + \frac{m\gamma}{4}})$ for some $\alpha > 0$, $\gamma > 1$ then (9) holds.

Here
$$D((-L)^{\alpha}) = \{z \in L_2(\Omega) : \sum_{k=1}^{\infty} \mu_k^{2\alpha} |\langle z, v_k \rangle|^2 < \infty\}.$$

$$\psi_{k_j} = \varphi_{k_j} E_{\beta}(-\lambda_{k_j} T^{\beta}) + f_{k_j} \int_0^T E_{\beta}(-\lambda_{k_j} (T - \tau)^{\beta}) q(\tau) d\tau, \ \ j \in \mathbb{N}.$$

$$\psi_{k_j} = arphi_{k_j} E_eta(-\lambda_{k_j} T^eta) + f_{k_j} \int_0^T E_eta(-\lambda_{k_j} (T- au)^eta) q(au) d au, \;\; j \in \mathbb{N}.$$

Transform them to the form

$$\frac{\lambda_{k_j}\psi_{k_j}}{\varphi_{k_i}} = \lambda_{k_j}E_{\beta}(-\lambda_{k_j}T^{\beta}) + \frac{f_{k_j}}{\varphi_{k_i}}\int_0^T \lambda_{k_j}E_{\beta}(-\lambda_{k_j}(T-\tau)^{\beta})q(\tau)d\tau, j \in \mathbb{N}$$

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The terms have the following behavior:

$$\underbrace{\frac{\lambda_{k_j}\psi_{k_j}}{\varphi_{k_j}}}_{\to M} = \underbrace{\lambda_{k_j}E_{\beta}(-\lambda_{k_j}T^{\beta})}_{\to \frac{1}{\Gamma(1-\beta)T^{\beta}}} + \underbrace{\frac{f_{k_j}}{\varphi_{k_j}}}_{\to 0} \underbrace{\int_0^T \lambda_{k_j}E_{\beta}(-\lambda_{k_j}(T-\tau)^{\beta})q(\tau)d\tau}_{\text{bounded}}.$$

$$\frac{1}{\Gamma(1-\beta)T^{\beta}} = M \quad \text{where} \quad M = \lim_{j \to \infty} \frac{\lambda_{k_j} \psi_{k_j}}{\varphi_{k_i}}. \tag{10}$$

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In case $T \geq e^{-\gamma_*} \approx$ 0.561, where $\gamma_* \approx$ 0.577 is the Euler-Mascheroni constant,

the function $\frac{1}{\Gamma(1-\beta)T^{\beta}}$ is strictly monotonic for $\beta \in (0,1)$.

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Treatment of IP3 (there λ and q are unknown).

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Again, assume that

there exists a subsequence k_j , $j \in \mathbb{N}$ such that

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Let us also recall that

$$\lambda_k = \lambda \mu_k$$

where

 μ_k is the eigenvalue of $-\Delta$ with Dirichlet boundary condition and λ_k is the eigenvalue of $-\lambda\Delta$ with Dirichlet boundary condition.

$$\psi_{k_j} = \varphi_{k_j} E_{\beta}(-\lambda \mu_{k_j} T^{\beta}) + f_{k_j} \int_0^T E_{\beta}(-\lambda \mu_{k_j} (T-\tau)^{\beta}) q(\tau) d\tau, \ \ j \in \mathbb{N}.$$

$$\psi_{k_j} = \varphi_{k_j} E_{\beta}(-\lambda \mu_{k_j} T^{\beta}) + f_{k_j} \int_0^T E_{\beta}(-\lambda \mu_{k_j} (T- au)^{\beta}) q(au) d au, \ \ j \in \mathbb{N}.$$

Transform them to the form

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and take the limit $j \to \infty$.

Terms of this equation have the following behavior:

$$\underbrace{\frac{\mu_{k_j}\psi_{k_j}}{\varphi_{k_j}}}_{\rightarrow M_1} = \underbrace{\mu_{k_j}E_{\beta}(-\lambda\mu_{k_j}T^{\beta})}_{\rightarrow \frac{1}{\lambda\Gamma(1-\beta)T^{\beta}}} + \underbrace{\frac{f_{k_j}}{\varphi_{k_j}}}_{\rightarrow 0} \underbrace{\int_0^T \mu_{k_j}E_{\beta}(-\lambda\mu_{k_j}(T-\tau)^{\beta})q(\tau)d\tau}_{\text{bounded}}$$



Therefore, we obtain

$$rac{1}{\lambda\Gamma(1-eta)T^eta}= extit{M}_1 \quad ext{where} \quad extit{M}_1=\lim_{j o\infty}rac{\mu_{k_j}\psi_{k_j}}{arphi_{k_j}}.$$

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This gives explicitly $\lambda = \frac{1}{M_1 \Gamma(1-\beta) T^{\beta}}$.

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Treatment of IP4 (there β , λ and q are unknown).

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Suppose that IP4 has a solution and consider the relations

$$\psi_{k_j} = \varphi_{k_j} E_{\beta}(-\lambda \mu_{k_j} T^{\beta}) + f_{k_j} \int_0^T E_{\beta}(-\lambda \mu_{k_j} (T- au)^{\beta}) q(au) d au, \ \ j \in \mathbb{N}.$$

Transform them to the form

$$\frac{\mu_{k_j}\psi_{k_j}}{\varphi_{k_i}} = \mu_{k_j}E_{\beta}(-\lambda\mu_{k_j}T^{\beta}) + \frac{f_{k_j}}{\varphi_{k_i}}\int_0^T \mu_{k_j}E_{\beta}(-\lambda\mu_{k_j}(T-\tau)^{\beta})q(\tau)d\tau, j \in \mathbb{N}$$

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We obtain the equation

$$\frac{1}{\lambda \Gamma(1-\beta) T^\beta} = \textit{M}_1 \quad \text{where} \quad \textit{M}_1 = \lim_{j \to \infty} \frac{\mu_{\textit{k}_j} \psi_{\textit{k}_j}}{\varphi_{\textit{k}_j}}.$$



Next we deduce the relation

$$\mu_{k_{j}}\left(\frac{\mu_{k_{j}}\psi_{k_{j}}}{\varphi_{k_{j}}}-M_{1}\right)=\mu_{k_{j}}\left(\mu_{k_{j}}E_{\beta}(-\lambda\mu_{k_{j}}T^{\beta})-\frac{1}{\lambda\Gamma(1-\beta)T^{\beta}}\right) + \frac{\mu_{k_{j}}f_{k_{j}}}{\varphi_{k_{j}}}\int_{0}^{T}\mu_{k_{j}}E_{\beta}(-\lambda\mu_{k_{j}}(T-\tau)^{\beta})q(\tau)d\tau, j \in \mathbb{N}.$$

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and take the limit $j \to \infty$.

Terms of this relation have the following behavior:

$$\underbrace{\frac{\mu_{k_{j}}\left(\frac{\mu_{k_{j}}\psi_{k_{j}}}{\varphi_{k_{j}}}-M_{1}\right)}_{\rightarrow M_{2}} = \underbrace{\mu_{k_{j}}\left(\mu_{k_{j}}E_{\beta}(-\lambda\mu_{k_{j}}T^{\beta})-\frac{1}{\lambda\Gamma(1-\beta)T^{\beta}}\right)}_{\rightarrow -\frac{1}{\lambda^{2}\Gamma(1-2\beta)T^{2\beta}}} + \underbrace{\frac{\mu_{k_{j}}f_{k_{j}}}{\varphi_{k_{j}}}\underbrace{\int_{0}^{T}\mu_{k_{j}}E_{\beta}(-\lambda\mu_{k_{j}}(T-\tau)^{\beta})q(\tau)d\tau}_{\text{bounded}}}_{\text{bounded}}$$

$$rac{1}{\lambda^2\Gamma(1-2eta)\mathcal{T}^{2eta}} = -M_2 \quad ext{where} \quad M_2 = \lim_{j o\infty} \mu_{k_j} \left(rac{\mu_{k_j}\psi_{k_j}}{arphi_{k_i}} - M_1
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We have the following system for β and λ :

$$\frac{1}{\lambda\Gamma(1-\beta)T^{\beta}}=M_{1},\quad \frac{1}{\lambda^{2}\Gamma(1-2\beta)T^{2\beta}}=-M_{2}.$$

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The function $\frac{(\Gamma(1-\beta))^2}{\Gamma(1-2\beta)}$ is strictly monotonic for $\beta \in (0,1)$.

Hence β is uniquely determined.



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Finally, the uniqueness of q can be shown as in case IP1.

Let us formulate theorems. Define

$$\mathcal{Q}=\{q:\, q(t)=0,\, t\in (T-\delta,T) \,\, ext{ for some } \,\, \delta\in (0,T), \,\,\, q\in L_2(0,T-\delta)\}$$
 and

 $\mathcal{F} = \{ f \in L_2(\Omega) : \text{ there exists a subsequence } k_i \text{ of } \mathbb{N} \text{ such that } f_{k_i} \neq 0, \ i \in \mathbb{N} \}.$

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Theorem 1

Let $f \in \mathcal{F}$ and $\varphi = \psi = 0$. If IP1 has a solution $q \in \mathcal{Q}$ then q = 0.

Theorem 2

Let $f \in \mathcal{F}$ and $T \ge e^{-\gamma_*}$, where γ_* is the Euler-Mascheroni constant. Moreover, assume that

$$\varphi \in L_2(\Omega)$$
, there exists a subsequence k_j , $j \in \mathbb{N}$ such that $\varphi_{k_j} \neq 0, \ j \in \mathbb{N}$, $\lim_{j \to \infty} \frac{f_{k_j}}{\varphi_{k_i}} = 0$.

If IP2 has solutions (β,q) , $(\tilde{\beta},\tilde{q})\in(0,1)\times\mathcal{Q}$ then $\beta=\tilde{\beta}$, $q=\tilde{q}$.

Theorem 3

Let $f \in \mathcal{F}$. Moreover, assume that

 $\varphi \in L_2(\Omega)$, there exists a subsequence k_j , $j \in \mathbb{N}$ such that $\varphi_{k_j} \neq 0, \ j \in \mathbb{N}$, $\lim_{j \to \infty} \frac{f_{k_j}}{\varphi_{k_i}} = 0$.

If $\operatorname{IP}3$ has solutions (λ,q) , $(\tilde{\lambda},\tilde{q})\in(0,\infty)\times\mathcal{Q}$ then $\lambda=\tilde{\lambda}$, $q=\tilde{q}$.

Theorem 4

Let $f \in \mathcal{F}$. Moreover, assume that

$$\varphi \in L_2(\Omega)$$
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If IP4 has solutions (β, λ, q) , $(\tilde{\beta}, \tilde{\lambda}, \tilde{q}) \in (0, 1) \times (0, \infty) \times \mathcal{Q}$ then $\beta = \tilde{\beta}$, $\lambda = \tilde{\lambda}$, $q = \tilde{q}$.

Thank you for the attention!