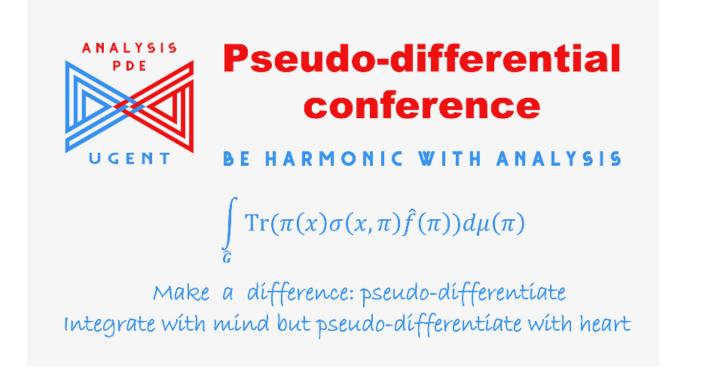
# The Harmonic Oscillator on The Heisenberg Group

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### Abstract

We present a notion of harmonic oscillator on the Heisenberg group  $\mathbf{H}_n$ . This operator forms a natural analogue of the harmonic oscillator on  $\mathbb{R}^n$ .

### Introduction

Our ansatz is based on a few reasonable assumptions: the harmonic oscillator on  $\mathbf{H}_n$  should be a negative sum of squares of operators related to the sub-Laplacian on  $\mathbf{H}_n$ , essentially self-adjoint with purely discrete spectrum, and its eigenvectors should be smooth functions and form an orthonormal basis of  $L^2(\mathbf{H}_n)$ . This leads to a differential operator on  $\mathbf{H}_n$  which is determined by the Dynin-Folland Lie algebra, a stratified 3-step nilpotent Lie algebra.

### Ansatz

Our approach is motivated by the following three realizations of the classical harmonic oscillator  $\mathcal{Q}_{\mathbb{R}^n}$ on  $\mathbb{R}^n$ :

- (R1) the negative sum of squares  $-\Delta + |x|^2$  of partial derivatives of order 1 and coordinate multiplication operators;
- (R2) the Weyl and Kohn-Nirenberg quantizations on  $\mathbb{R}^n$  of the symbol  $\sigma(x,\xi) := |\xi|^2 + |x|^2$  with
- (R3) the image  $d\rho_1(-\mathcal{L}_{\mathbf{H}_n})$  of the negative sub-Laplacian  $-\mathcal{L}_{\mathbf{H}_n}$  on  $\mathbf{H}_n$  under the infinitesimal Schrödinger representation  $d\rho_1$  (of Planck's constant equal to 1) of the Heisenberg Lie algebra  $\mathfrak{h}_n$ .

The Schrödinger representation  $\rho_1$  of  $\mathbf{H}_n$  acting on  $L^2(\mathbb{R}^n)$  and the associated Lie algebra representation, naturally acting on  $\mathcal{S}(\mathbb{R}^n)$ , clearly relate each of the realisations (R1) - (R3) to the others. One can expect that similar realisations should be available for the canonical harmonic oscillator on  $\mathbf{H}_n$ .

### Main Result

The harmonic oscillator  $\mathcal{Q}_{\mathbf{H}_n}$  on the Heisenberg group  $\mathbf{H}_n$  has a purely discrete spectrum spec $(\mathcal{Q}_{\mathbf{H}_n}) \subset$  $(0,\infty)$ . The number of its eigenvalues, counted with multiplicaties, which are less or equal to  $\lambda>0$  is asymptotically (as  $\lambda \to \infty$ ) given by

$$N(\lambda) \sim \lambda^{rac{6n+3}{2}},$$

and the magnitude of the eigenvalues is asymptotically equal to

$$\lambda_k \sim k^{\frac{2}{6n+3}}$$
 for  $k=1,2,\ldots$ 

Moreover, the eigenvectors of  $\mathcal{Q}_{\mathbf{H}_n}$  are in  $\mathcal{S}(\mathbf{H}_n)$  and form an orthonormal basis of  $L^2(\mathbf{H}_n)$ .

$$Q_{H_{4}} = -\int_{H_{4}} + \chi_{3}^{2} = -(2\chi_{4}^{2} + 2\chi_{2}^{2}) - \frac{1}{4}(\chi_{4}^{2} + \chi_{2}^{2})2\chi_{3}^{2} + (\chi_{4} 2\chi_{2} - \chi_{2} 2\chi_{4})2\chi_{3} + \chi_{3}^{2}$$

Figure 1: The Harmonic Oscillator on  $H_1$ .

### Definition

The Dynin-Folland Lie group  $\mathbf{H}_{n,2} = \mathbb{R}^{2n+2} \rtimes \mathbf{H}_n$  acts on  $f \in L^2(\mathbf{H}_n)$  by the unitary irreducible representation  $(\pi(z,y,x)f)(t) = e^{iz}e^{i\langle t\cdot \frac{1}{2}x,y\rangle}f(t\cdot x),$ 

where  $t \cdot \frac{1}{2}x$  and  $t \cdot x$  denote the  $\mathbf{H}_n$ -group products of the corresponding coordinate vectors. For the basis  $\{X_1,\ldots,Y_{2n+1},Z\}$  of its Lie algebra  $\mathfrak{h}_{n,2}$  we define the **harmonic oscillator** on  $\mathbf{H}_n$  by

$$\mathcal{Q}_{\mathbf{H}_n} := d\pi (-\mathcal{L}_{\mathbf{H}_{n,2}}) = -d\pi (X_1)^2 - \ldots - d\pi (X_{2n})^2 - d\pi (Y_{2n+1})^2,$$

where  $-\mathcal{L}_{\mathbf{H}_n}$  is the sub-Laplacian on  $\mathbf{H}_{n,2}$ . Its natural domain includes the smooth vectors  $\mathcal{H}_{\pi}^{\infty} \cong \mathcal{S}(\mathbf{H}_n)$ .

## Interpretation

The essentially self-adjoint differential operator  $\mathcal{Q}_{\mathbf{H}_n}$ on  $\mathbf{H}_n$  admits analogues of (R1) – (R3):

- (R1') the differential operator  $-\mathcal{L}_{\mathbf{H}_n} + x_{2n+1}^2$ ;
- (R2') i) the Dynin-Weyl quantization on  $\mathbf{H}_n$  of the symbol  $\sigma(x,\xi) := \xi_1^2 + \dots + \xi_{2n}^2 + x_{2n+1}^2$  with  $(x,\xi) \in \mathbb{R}^{2n+1} \times \widehat{\mathbb{R}}^{2n+1};$ 
  - ii) the Kohn-Nirenberg quantization, in the sense of [3], of the operator-valued symbol on  $\mathbf{H}_n \times \hat{\mathbf{H}}_n$  $\sigma(x, \rho_{\lambda}) := -\rho_{\lambda}(X_1)^2 - \ldots - \rho_{\lambda}(X_{2n})^2 + x_{2n+1}^2;$
- (R3') the image  $d\pi(-\mathcal{L}_{\mathbf{H}_{n,2}})$  of the sub-Laplacian  $\mathcal{L}_{\mathbf{H}_{n,2}}$ on  $\mathbf{H}_{n,2}$ , here by definition.

#### Methods

 $\mathcal{Q}_{\mathbf{H}_n}$  has purely discrete spectrum in  $(0, \infty)$  by [5]. The asymptotic growth rate of its eigenvalues is obtained via a powerful method developed in [8]. The number of eigenvalues (with multiplicities), asymptotically behaves like the volumes of certain subsets of the coadjoint orbit  $\mathcal{O}_{\pi} \subset \mathfrak{h}_{n,2}^*$  corresponding to representation  $\pi \in \hat{\mathbf{H}}_n$ . The subsets in question are determined (up to a multiplicative constant) by a (any) homogeneous quasi-norm on  $\mathfrak{g}^*$ . The flatness of  $\mathcal{O}_{\pi}$  and a convenient choice of quasi-norm facilitate the computations substantially.

### Remark

The power  $\frac{6n+3}{2}$  intimately related to the homogeneous structure of  $\mathfrak{h}_{n,2}$ : the nominator 6n+3 is the homogenous dimension of the first two strata  $\mathfrak{g}_1 \oplus \mathfrak{g}_2 \subseteq \mathfrak{h}_{n,2}$ , while the denominator 2 is the homogeneous degree of  $-\mathcal{L}_{\mathbf{H}_{n}}$ .

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