

Fractional Diffusion Equations

Kai Diethelm



Seasonal School 2025

—

Ghent Analysis & PDE Center
13–24 October 2025

- ① Classical Diffusion
- ② Generalization of the Classical Diffusion Model
- ③ Time-Fractional Diffusion Equations
- ④ Space-Fractional and Space-and-Time-Fractional Diffusion Equations

Random Walk in 1D

- at time $t = 0$, random walker is in position $x = 0$
- random walker jumps at times $t = n\tau$ ($n = 1, 2, \dots$)
where $\tau > 0$ is a fixed constant
- jump length = $\xi > 0$ (the **lattice constant**; also fixed)
- direction of jump is randomly chosen;
equal probabilities for jumps to right and to left
- Notation: $P(x, t) \cdot \xi =$ probability that random walker's position
at time t is in interval $[x, x + \xi)$
- **master equation**:

$$P(x, t + \tau) = \frac{1}{2}P(x + \xi, t) + \frac{1}{2}P(x - \xi, t)$$

Master equation:

$$P(x, t + \tau) = \frac{1}{2}P(x + \xi, t) + \frac{1}{2}P(x - \xi, t)$$

Taylor expansion:

$$P(x, t + \tau) = P(x, t) + \tau \frac{\partial P}{\partial t}(x, t) + O(\tau^2)$$

$$P(x \pm \xi, t) = P(x, t) \pm \xi \frac{\partial P}{\partial x}(x, t) + \frac{\xi^2}{2} \frac{\partial^2 P}{\partial x^2}(x, t) + O(\xi^3)$$

Plug into master equation and rearrange:

$$\frac{\partial P}{\partial t}(x, t) = \frac{\xi^2}{2\tau} \frac{\partial^2 P}{\partial x^2}(x, t) + \xi \cdot O\left(\frac{\xi^2}{\tau}\right) + O(\tau)$$

$$\frac{\partial P}{\partial t}(x, t) = \frac{\xi^2}{2\tau} \frac{\partial^2 P}{\partial x^2}(x, t) + \xi \cdot O\left(\frac{\xi^2}{\tau}\right) + O(\tau)$$

Continuum limit ($\tau \rightarrow 0$ and $\xi \rightarrow 0$ coupled such that **diffusion coefficient** $d_1 = \lim_{\tau \rightarrow 0, \xi \rightarrow 0} \xi^2/(2\tau)$ exists):

$$\frac{\partial P}{\partial t}(x, t) = d_1 \frac{\partial^2 P}{\partial x^2}(x, t) \quad (\text{classical diffusion equation})$$

Can be generalized to n -dimensional case:

$$\frac{\partial P}{\partial t}(x, t) = d_1 \Delta P(x, t) \quad \text{where} \quad \Delta = \sum_{j=1}^n \frac{\partial^2}{\partial x_j^2} \quad (\text{Laplace operator})$$

- Diffusion equation:

$$\frac{\partial P}{\partial t}(x, t) = d_1 \frac{\partial^2 P}{\partial x^2}(x, t)$$

- Random walker located in position $x = 0$ at time $t = 0$
 \Rightarrow initial condition $P(x, 0) = \delta(x)$ (Dirac distribution)
- Solution of initial value problem:

$$P(x, t) = \frac{1}{\sqrt{4\pi d_1 t}} \exp\left(-\frac{x^2}{4d_1 t}\right)$$

- For fixed t , $P(\cdot, t)$ is (spatial) Gaussian distribution with mean value $\mu = 0$ and standard deviation $\sigma(t) = \sqrt{2d_1 t}$
 \Rightarrow mean squared displacement of random walker at time t is

$$\sigma^2(t) = 2d_1 t$$

Classification of diffusion processes according to behaviour of mean squared displacement at time t (when t is large):

$$\sigma^2(t) \sim t^\alpha \quad \text{for } t \rightarrow \infty$$

- classical diffusion (see above): $\alpha = 1$
- subdiffusion: $\alpha < 1$
- superdiffusion: $\alpha > 1$
- ballistic diffusion: $\alpha = 2$

Next steps:

Find mathematical models that exhibit such kinds of behaviour

Random Walk in 1D

- at time $t = 0$, random walker is in position $x = 0$
- random walker jumps at times $t = n\tau$ ($n = 1, 2, \dots$)
where $\tau > 0$ is a fixed constant
- jump length $= \xi > 0$ (the lattice constant also fixed)
- direction of jump is randomly chosen;
equal probabilities for jumps to right and to left

Many deterministic assumptions!

- at time $t = 0$, random walker is in position $x = 0$
- joint probability distribution function $\psi(x, t)$ (so-called **jump PDF**) governs **lengths of jumps** and **waiting time between successive jumps**
- jump length PDF:

$$\lambda(x) = \int_0^{\infty} \psi(x, t) dt$$

- waiting time PDF:

$$w(t) = \int_{-\infty}^{\infty} \psi(x, t) dx$$

- main characteristic features of CTRW model:

- ▶ characteristic waiting time $T = \int_0^{\infty} tw(t) dt$ (first moment of w)
- ▶ jump length variance $\Sigma^2 = \int_{-\infty}^{\infty} x^2 \lambda(x) dx$

- characteristic waiting time $T = \int_0^{\infty} tw(t) dt$
- jump length variance $\Sigma^2 = \int_{-\infty}^{\infty} x^2 \lambda(x) dx$

	$T < \infty$	$T = \infty$
$\Sigma^2 < \infty$	classical diffusion (Brownian motion)	(non-Markovian) subdiffusion (time-fractional PDE)
$\Sigma^2 = \infty$	Markovian Lévy flights (space-fractional PDE)	non-Markovian Lévy flights (space-and-time-fractional PDE)

Simplifying assumption:

- jump lengths and waiting times are independent random variables
- $\psi(x, t) = \lambda(x)w(t)$

$$\begin{aligned}\eta(x, t) &= \text{PDF of just having arrived in position } x \text{ at time } t \\ &= \int_{-\infty}^{\infty} \int_0^{\infty} \eta(\xi, \tau) \psi(x - \xi, t - \tau) d\tau d\xi + P_0(x) \delta(t) \\ & (P_0 = \text{initial condition})\end{aligned}$$

Then,

$$\begin{aligned}P(x, t) &= \text{PDF of being in position } x \text{ at time } t \\ &= \int_0^t \eta(x, \tau) \Psi(t - \tau) d\tau\end{aligned}$$

where $\Psi(t) = 1 - \int_0^t w(\tau) d\tau$ (“no jump in time interval $[0, t]$ ”)

Application of Fourier transform in space and Laplace transform in time to this PDF yields

$$\begin{aligned}\mathcal{FL}P(\omega, s) &= \frac{1 - \mathcal{L}w(s)}{s} \cdot \frac{\mathcal{F}P_0(\omega)}{1 - \mathcal{FL}\psi(\omega, s)} \\ &= \frac{1 - \mathcal{L}w(s)}{s} \cdot \frac{\mathcal{F}P_0(\omega)}{1 - \mathcal{F}\lambda(\omega)\mathcal{L}w(s)}\end{aligned}\quad (*)$$

Special case: waiting time PDF with **heavy tail**

$$w(t) \sim A_\alpha \left(\frac{\tau}{t}\right)^{\alpha+1} \quad \text{for } t \rightarrow \infty$$

- $\alpha \in (0, 1)$
- $\tau > 0$: characteristic time constant
- $A_\alpha > 0$: normalization factor

Properties:

- w is integrable,
characteristic waiting time (first moment of w): $T = \infty$.
- Tauberian theorem: $\mathcal{L}w(s) \sim 1 - (s\tau)^\alpha$ as $s \rightarrow 0$

Concrete example:

$$\text{waiting time PDF } w(t) = -\frac{d}{dt}E_{\alpha,1}(-t^\alpha)$$

$$\text{with Mittag-Leffler function } E_{\alpha,\beta}(z) = \sum_{k=0}^{\infty} \frac{z^k}{\Gamma(\alpha k + \beta)} \text{ and } 0 < \alpha < 1$$

$$\text{Gaussian jump length PDF } \lambda(x) = \frac{1}{\sqrt{4\pi\sigma^2}} \exp\left(-\frac{x^2}{4\sigma^2}\right)$$

$$\text{such that } \Sigma^2 = 2\sigma^2 < \infty$$

$$\mathcal{L}w(s) = \frac{1}{1+s^\alpha}, \quad \mathcal{F}\lambda(\omega) \sim 1 - \sigma^2\omega^2 \quad (\omega \rightarrow 0)$$

Recalling

$$\mathcal{F}\mathcal{L}P(\omega, s) = \frac{1 - \mathcal{L}w(s)}{s} \cdot \frac{\mathcal{F}P_0(\omega)}{1 - \mathcal{F}\lambda(\omega)\mathcal{L}w(s)} \quad (*)$$

and plugging in the above statements $\mathcal{L}w(s) = \frac{1}{1+s^\alpha}$ and $\mathcal{F}\lambda(\omega) \sim 1 - \sigma^2\omega^2$, we obtain for $s \rightarrow 0$ and $\omega \rightarrow 0$ that

$$\mathcal{F}\mathcal{L}P(\omega, s) \sim \frac{\mathcal{F}P_0(\omega)/s}{1 + d_\alpha\omega^2s^{-\alpha}}$$

with some constant d_α , or

$$(1 + d_\alpha\omega^2s^{-\alpha})\mathcal{F}\mathcal{L}P(\omega, s) \sim \frac{\mathcal{F}P_0(\omega)}{s}$$

Exploiting Fourier and Laplace transforms and their properties, this yields for large t and $|x|$

$$P(x, t) - P_0(x) = d_\alpha J_{0,t}^\alpha P_{xx}(x, t)$$

where $J_{0,t}^\alpha$ is the Riemann-Liouville integral of order α w.r.t. t .

Equivalent formulation: IVP for **time-fractional diffusion equation**

$$D_{*0,t}^\alpha P(x, t) = d_\alpha \frac{\partial^2 P}{\partial x^2}(x, t), \quad P(x, 0) = P_0(x),$$

where $D_{*0,t}^\alpha$ is the Caputo derivative of order α w.r.t. t .

Laplace transform techniques lead to

$$\text{mean squared displacement} = \sigma^2(t) = \frac{2d_\alpha}{\Gamma(1+\alpha)} t^\alpha$$

$$D_{*0,t}^{\alpha} u(x, t) = -Lu(x, t) + F(x, t), \quad x \in G \subset \mathbb{R}^n, \quad t \in (0, T) \quad (\square)$$

- $\alpha \in (0, 1]$; $D_{*0,t}^{\alpha}$ is the Caputo derivative of order α w.r.t. t
- G is open and bounded domain in \mathbb{R}^n
- $Lu(x, t) = -\operatorname{div}(p(x) \operatorname{grad} u(x, t)) + q(x)u(x, t)$
- $p \in C(\bar{G})$, $q \in C(\bar{G})$, $p(x) > 0$ and $q(x) \geq 0$ for $x \in \bar{G}$

– L is elliptic differential operator of order 2,

$$\begin{aligned} -Lu(x, t) &= \sum_{k=1}^n \left(p(x) \frac{\partial^2 u}{\partial x_k^2}(x, t) + \frac{\partial p}{\partial x_k}(x) \frac{\partial u}{\partial x_k}(x, t) \right) - q(x)u(x, t) \\ &= p(x) \Delta u(x, t) + (\operatorname{grad} p(x, t), \operatorname{grad} u(x, t)) - q(x)u(x, t) \end{aligned}$$

Problem (\square) reduces to classical parabolic PDE when $\alpha = 1$

Same conditions as in classical parabolic PDE
(with time derivative of order 1):

$$\begin{aligned} u(x, 0) &= u_0(x) && \text{for all } x \in \bar{G}, \\ u(x, t) &= v(x, t) && \text{for all } x \in \partial G \text{ and } t \in [0, T] \end{aligned}$$

Definition

A function $u : \bar{G} \times [0, T]$ is called a **solution** to the IBVP above if

- u is continuous in its domain of definition,
- u is twice differentiable with respect to x in G ,
- u is once differentiable with respect to $t \in (0, T]$ such that the first time derivative is in $L^1(0, T)$,
- u satisfies the differential equation (\square) and the initial and boundary conditions above.

General assumptions (can be weakened):

- forcing function: $F \in C(\bar{G} \times [0, T])$
- initial function: $u_0 \in C(\bar{G})$
- boundary function: $v \in C(\partial G \times [0, T])$

Theorem

Under the given assumptions, the time-fractional IBVP above possesses at most one solution.

The solution depends continuously on F , u_0 and v in the following sense: If u_1 and u_2 are solutions to the PDE (\square) with forcing, initial and boundary functions F_1 , u_{01} and v_1 or F_2 , u_{02} and v_2 , respectively, then

$$\begin{aligned} \|u_1 - u_2\|_{L_\infty(\bar{G} \times [0, T])} \leq & \max \left\{ \|u_{01} - u_{02}\|_{L_\infty(\bar{G})}, \|v_1 - v_2\|_{L_\infty(\partial G \times [0, T])} \right\} \\ & + \frac{T^\alpha}{\Gamma(1 + \alpha)} \|F_1 - F_2\|_{L_\infty(\bar{G} \times [0, T])} \end{aligned}$$

Theorem

Let u be a solution to the IBVP above under the given assumptions and subject to the condition $F(x, t) \leq 0$ for all x and t . Then

- *either $u(x, t) \leq 0$ for all x and t*
- *or $u(x, t)$ attains its positive maximum on the part $S = (\bar{G} \times \{0\}) \cup (\partial G \times [0, T])$ of the boundary of its domain of definition,*

i.e.

$$u(x, t) \leq \max \left\{ 0, \max_{(\xi, \tau) \in S} u(\xi, \tau) \right\} \quad \text{for all } x \text{ and } t.$$

Homogeneous PDE and BC, general IC:

$$D_{*0,t}^{\alpha} u(x, t) = -Lu(x, t)$$

for all $x \in G$ and $t \in (0, T)$,

$$u(x, 0) = u_0(x)$$

for all $x \in \bar{G}$,

$$u(x, t) = 0$$

for all $x \in \partial G$ and $t \in [0, T]$

ansatz (Fourier):

$$u(x, t) = X(x)T(t)$$

separation of variables:

$$\frac{D_{*0}^{\alpha} T(t)}{T(t)} = -\frac{LX(x)}{X(x)} = -\lambda$$

Decomposition:

solution: $T(t) = E_{\alpha,1}(-\lambda t^{\alpha})$

$$D_{*0}^{\alpha} T(t) = -\lambda T(t),$$

$$L(X)(x) = -\lambda X(x), \quad X(x) = 0 \text{ for } x \in \partial G.$$

classical eigenvalue problem for elliptic operator L

Consequences:

- eigenvalue problem ($L + BC$) is positive definite and self-adjoint
- eigenvalue problem has countably many eigenvalues λ_j , $j \in \mathbb{N}$, all of which are in $(0, \infty)$ and have finite multiplicity (WLOG $0 < \lambda_1 \leq \lambda_2 \leq \lambda_3 \leq \dots$)
- denote the eigenfunction corresponding to λ_j by X_j
- any function of the form $u_j(x, t) = c_j E_{\alpha,1}(-\lambda_j t^\alpha) X_j(x)$ satisfies the PDE and the boundary condition

Theorem

If ∂G is smooth then the function

$$u(x, t) = \sum_{j=1}^{\infty} (u_0, X_j) E_{\alpha,1}(-\lambda_j t^\alpha) X_j(x)$$

is a formal solution to the fractional initial-boundary value problem. Under suitable assumptions, it is even an actual solution.

assumption: compatibility of IC & BC; smoothness of given data

Classical parabolic IBVP ($\alpha = 1$):

infinitely differentiable w.r.t. t on $[0, T]$

$$u(x, t) = \sum_{j=1}^{\infty} (u_0, X_j) \exp(-\lambda_j t) X_j(x)$$

exponential decay as $t \rightarrow \infty$

Fractional parabolic IBVP ($0 < \alpha < 1$):

infinitely differentiable w.r.t. t on $(0, T]$; not differentiable at 0

$$u(x, t) = \sum_{j=1}^{\infty} (u_0, X_j) E_{\alpha, 1}(-\lambda_j t^{\alpha}) X_j(x)$$

algebraic decay $\sim t^{-\alpha}$ as $t \rightarrow \infty$

- Behaviour with respect to x (smoothness etc.): same in both cases
- Behaviour with respect to t : remember $\lambda_j > 0$ for all j

References

- D. Baleanu, K. Diethelm, E. Scalas & J. J. Trujillo (2017), DOI 10.1142/10044
- Y. Luchko (2015), DOI 10.1007/978-3-642-54551-1_60
- M. Meerschaert & A. Sikorskii (2019), DOI 10.1515/9783110560244
- R. Metzler & J. Klafter (2000), DOI 10.1016/S0370-1573(00)00070-3

Thank you for your attention!

Contact: kai.diethelm@thws.de